

Local and Global Measures of the Stress Distribution in Abrupt Contraction-Expansions

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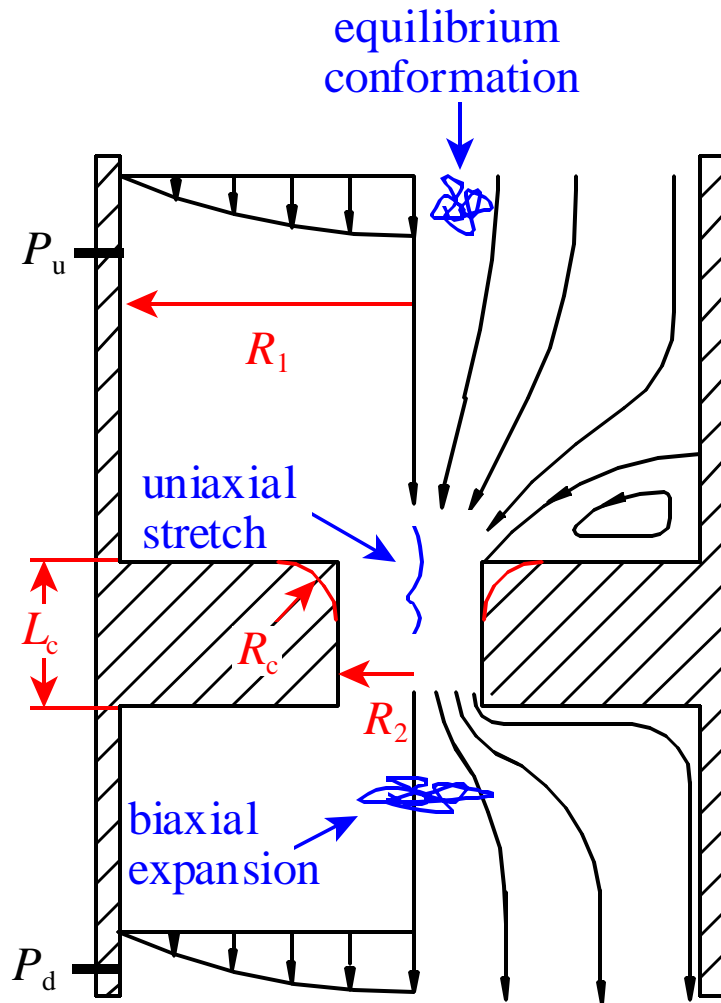


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The Axisymmetric Contraction-Expansion



(Cartalos & Piau, 1992; Szabo *et al.*, 1997; Rothstein & McKinley, 1999)



- Ⓒ Upstream contraction flow is numerical and experimental benchmark problem
- Ⓒ Complex flow containing mixture of shear near walls and extensional effects near contraction
- Ⓒ Combination of global and local flow measurements used to extensively characterize flow

- Ⓒ Characteristic Deborah number

$$De = l \dot{\gamma} = l Q / \pi R_2^3$$

- Ⓒ Creeping flow regime

$$Re < 10^{-3}$$

$$R_1 / R_2 = \beta = 2, 4, 8$$

$$L_c / R_2 = 0.5, 1, 2$$

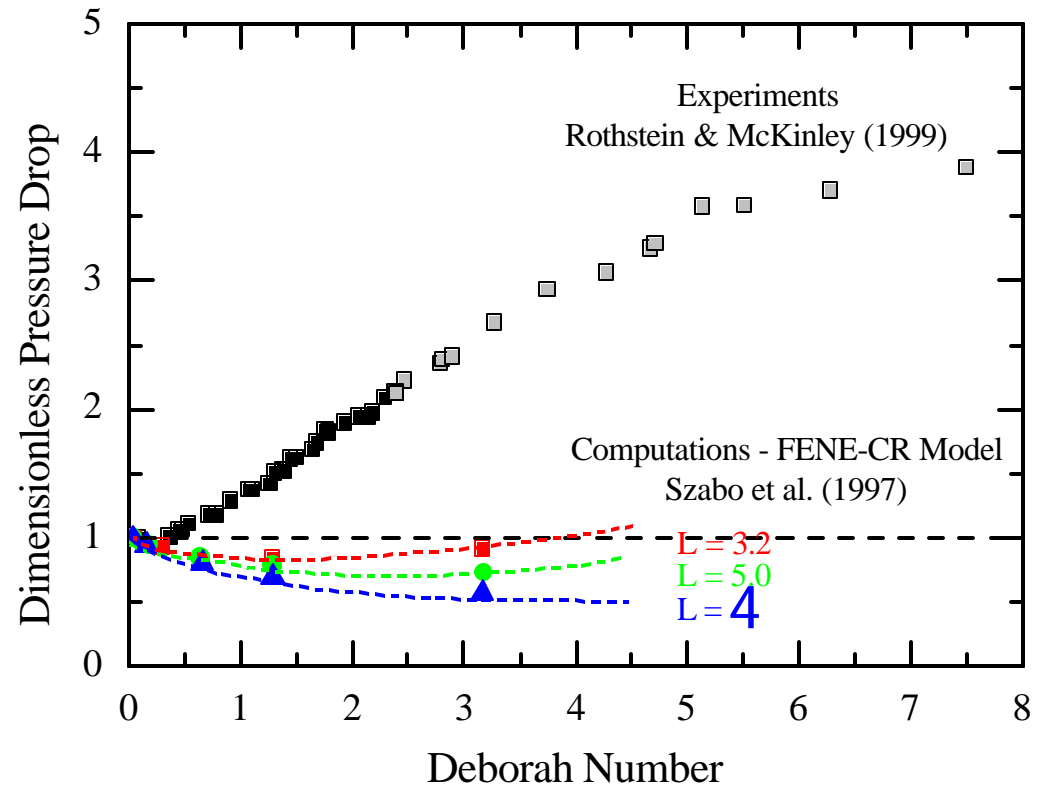
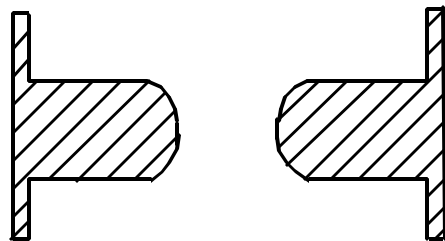
$$R_c = 0, 0.1R_2, 0.2R_2, 0.5R_2$$

Motivation



C There is a lack of agreement between experiments and computations in non-homogeneous flows

4:1:4 Contraction-Expansion
with Rounded Corners



C These discrepancies may be due to the internal, purely-dissipative stresses arising from non-equilibrium molecular conformations.

C Presence of stress-conformation hysteresis in transient uniaxial elongation suggests need for *global* (pressure drop) and *local* (birefringence) probes of conformation and stress.

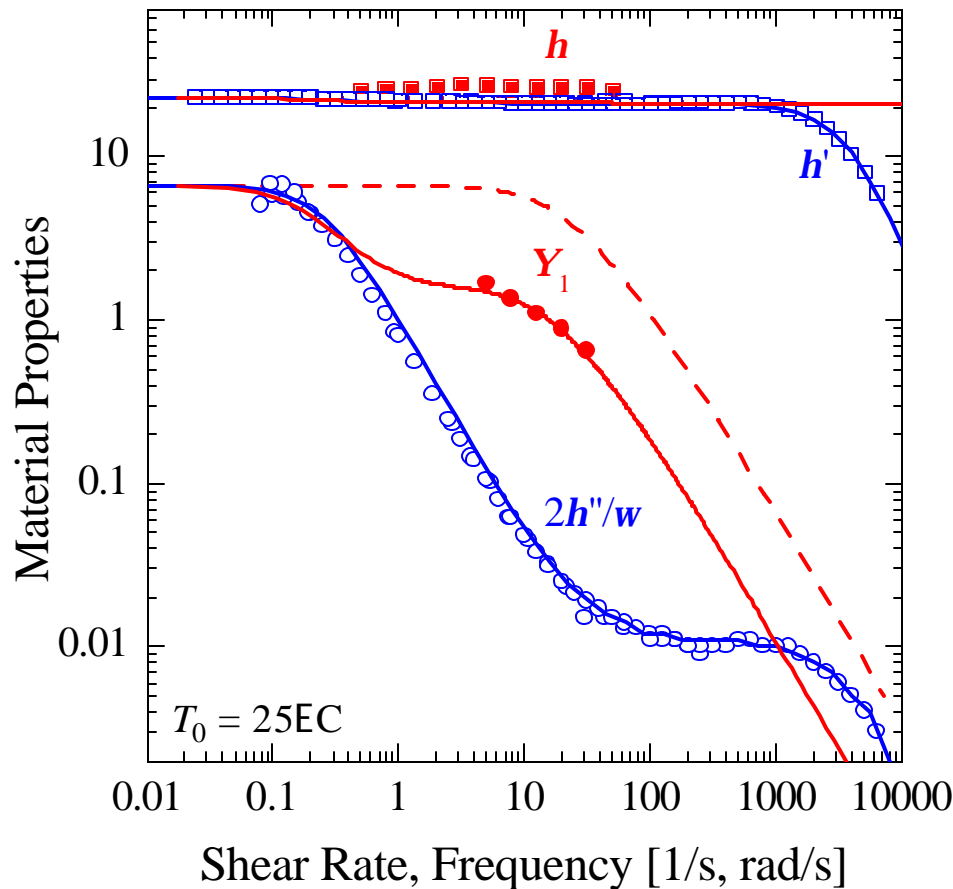
Fluid Shear Rheology



Monodisperse polystyrene $M_w = 2.03 \times 10^6$ g/mol dissolved in oligomeric polystyrene

Dilute solution with concentration $c = 0.025$ wt% γ $c/c^* = 0.23$

Model viscoelastic fluid to probe flow in the absence of polydispersity and inertial effects



Small amplitude oscillatory shear rheology:
 well fit by Rouse-Zimm bead-spring model

$h^* \cdot 0.1$ γ dominant hydrodynamic interactions

$$\delta_1 = 3.08\text{s}$$

$$\delta_i = \delta_1 / i^{1.77}$$

$$O_s/O_0 = 0.92$$

Steady shear rheology:

poorly fit by FENE-P model '&&' ($L = 88$)

well fit by Bird-DeAguiar model '&' ($L = 88, F = 0.62, \beta = 1.0$)

Extensional rheology also **very** important

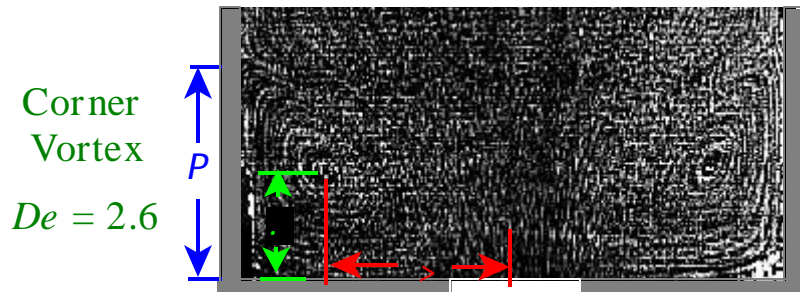
Vortex Growth Dynamics

C At moderate Deborah numbers, **two distinct patterns of vortex growth** exist

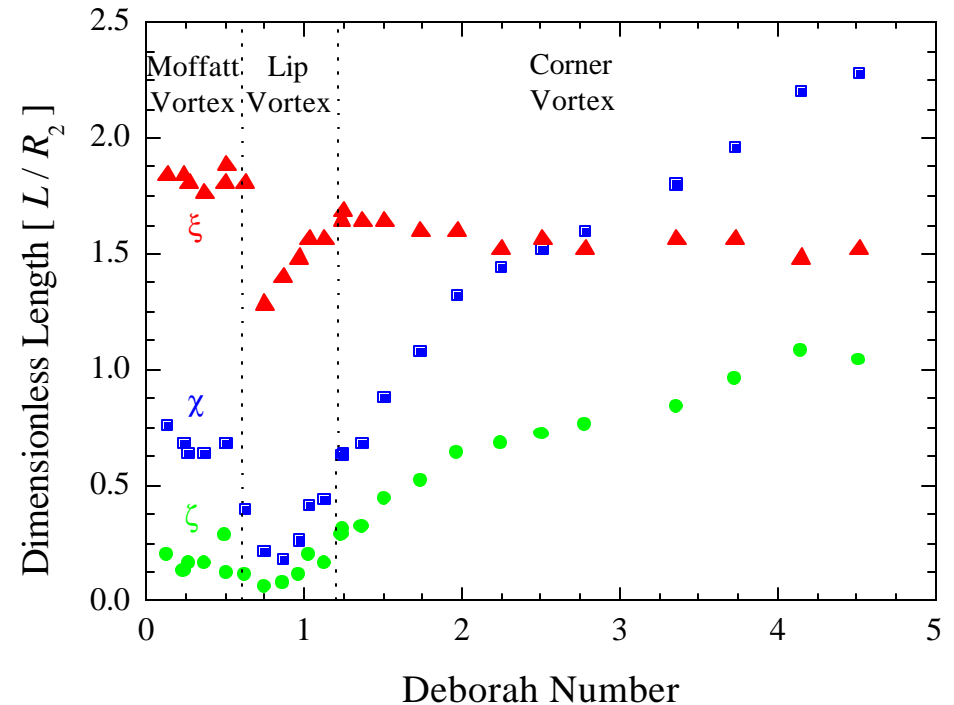
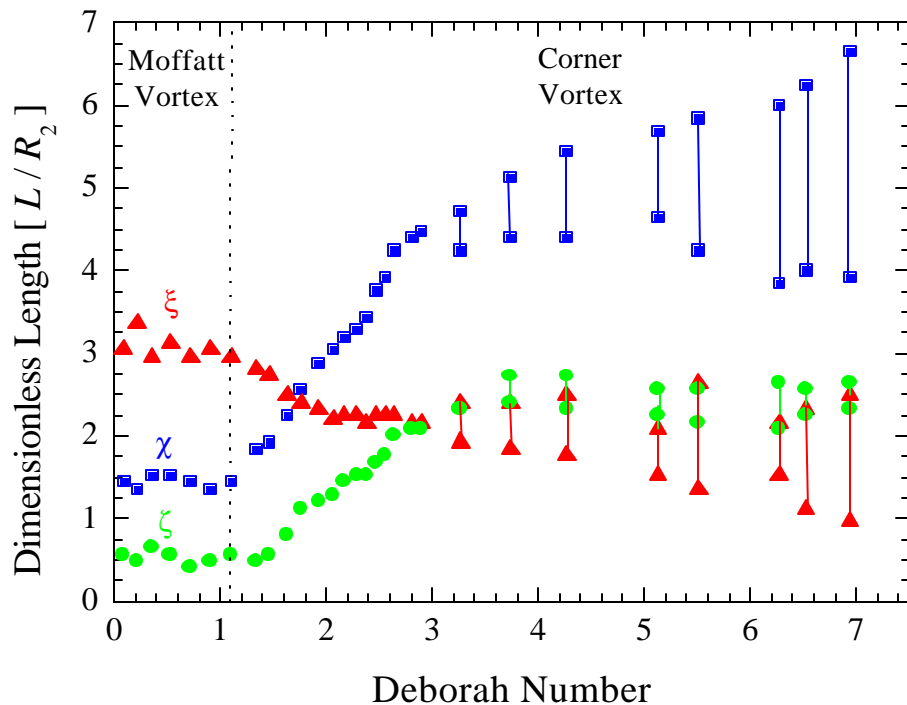
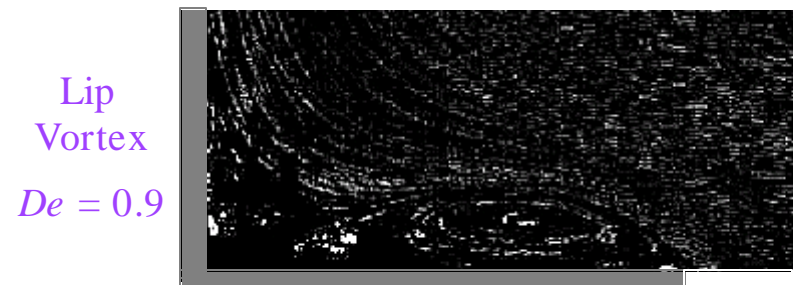
C For $b \geq 4$, vortex grows out from salient corner, grows upstream and eventually becomes unstable

C For $b = 2$, 'lip' vortex emerges near re-entrant corner, grows toward salient corner and upstream

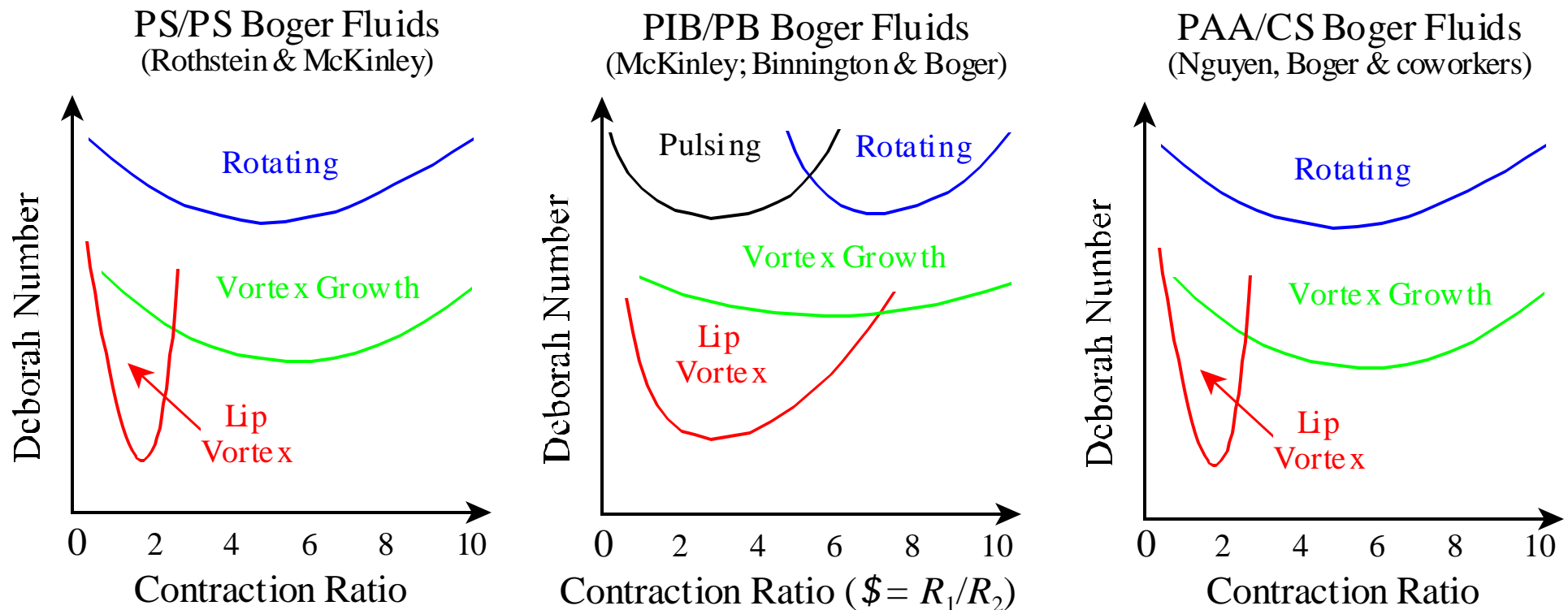
4:1:4 Contraction/Expansion



2:1:2 Contraction/Expansion



Flow Stability Diagram for PS/PS Boger Fluid



C Flow stability diagram of PS/PS Boger fluid is very similar to PAA/CS Boger fluid and dissimilar to PIB/PB Boger fluid

C Why do vortex growth dynamics depend on both contraction ratio and test fluid used?

Y Many have speculated that fluid dependence of flow structure arises from differences in extensional viscosity of different Boger fluids.

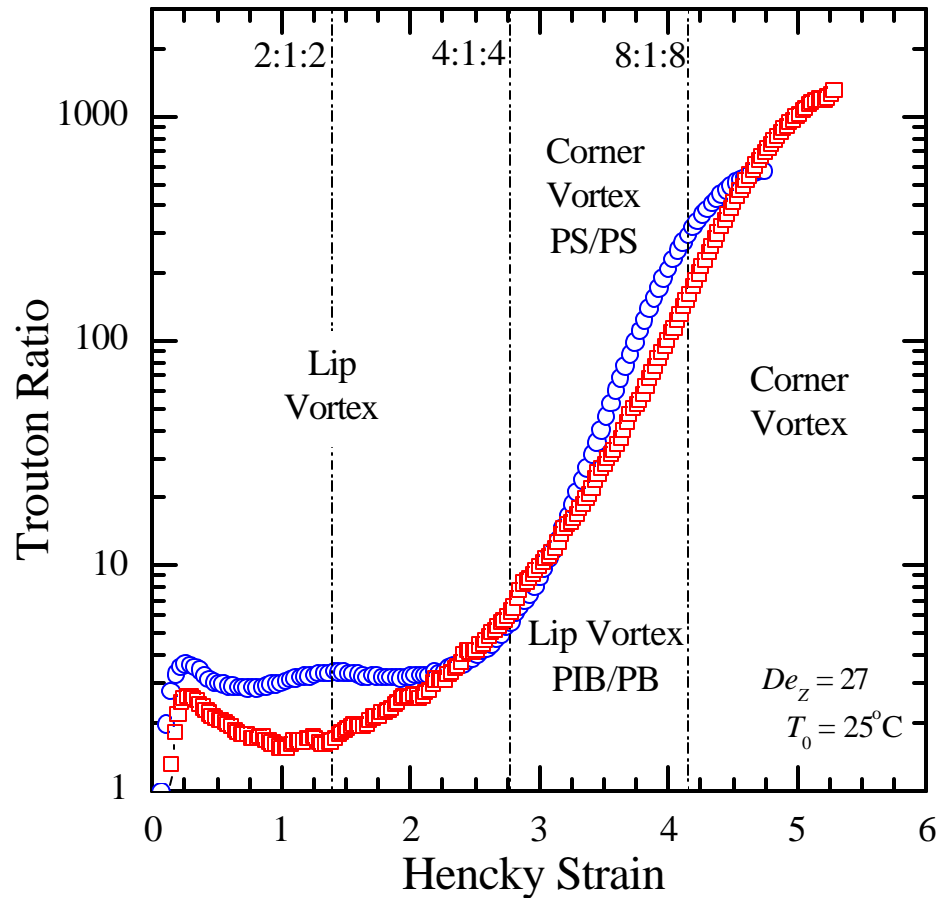
Transient Extensional Rheology

- C **PIB/PB** and **PS/PS** Boger fluids exhibit similar stress growth
- C Results are approximately independent of $De_z = \lambda_z \dot{\epsilon}$ (for $De_z > 1$)

Accumulated strain along centerline

$$\epsilon = \int_0^{t_1} \dot{\epsilon} dt = \int_{v_z(z=-\infty)}^{v_z(z=0.5L_c)} \frac{dv_z}{v_z} = 2 \ln \beta$$

- C Lip vortex present for
 - PS/PS $\beta < 2.77 \ \forall Tr = (\tau_{zz} - \tau_{rr}) / \eta_0 \dot{\epsilon} < 6$
 - PIB/PB $\beta < 4.16 \ \forall Tr < 200$



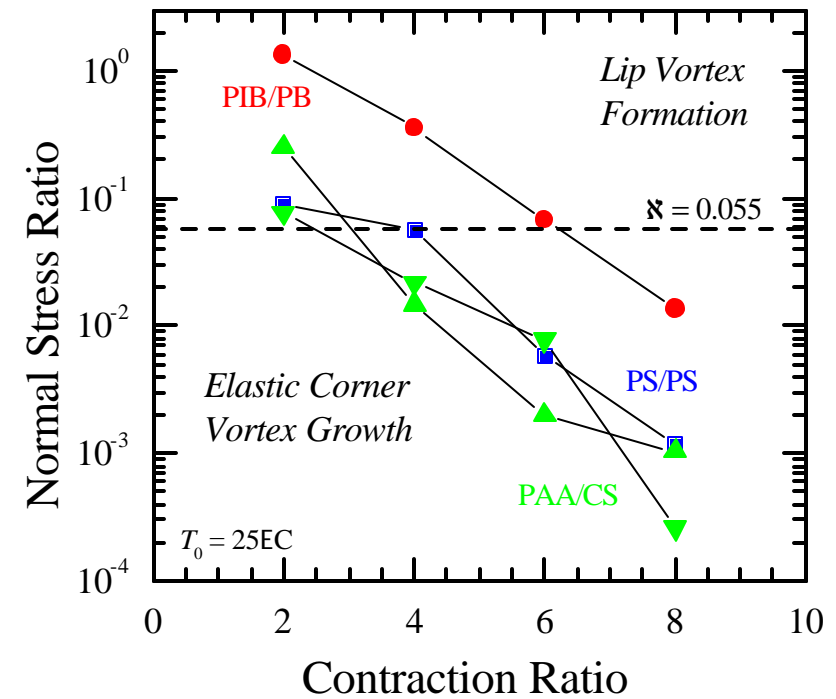
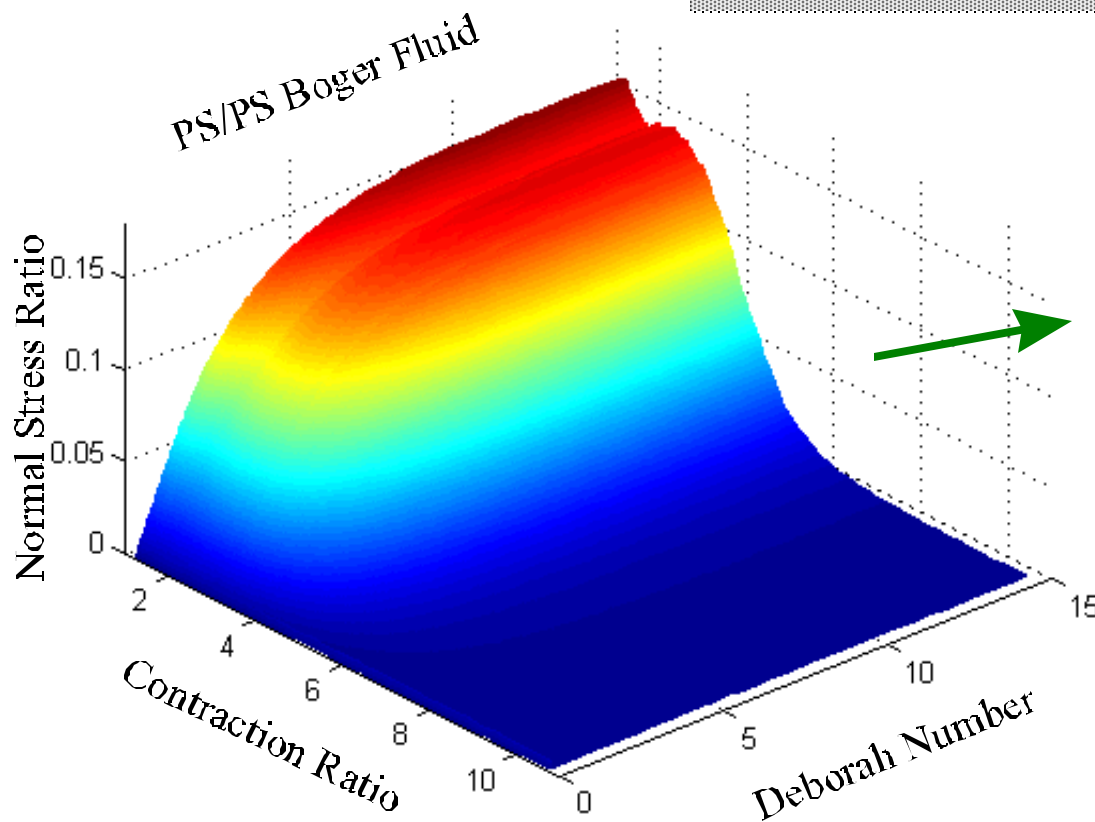
- C Independent knowledge of extensional and shear rheology alone does not explain choice of lip or corner vortex!
- C Need to understand how upstream shear flow affects extensional flow entering contraction.

Normal Stress Ratio



Systematic differences between fluids can be understood by considering relative importance of elastic normal stresses generated in shear to elastic normal stresses generated in transient uniaxial extension.

$$\mathfrak{N} = \frac{N_1 / \eta_0 \dot{\gamma}}{(\tau_{zz} - \tau_{rr}) / \eta_0 \dot{\epsilon}} = \frac{S_R(\dot{\gamma})}{T_R(\dot{\epsilon})}$$



For each fluid, a lip vortex does not develop for normal stress ratios ($\mathfrak{N} < 0.055$).

Need for Birefringence Measurements



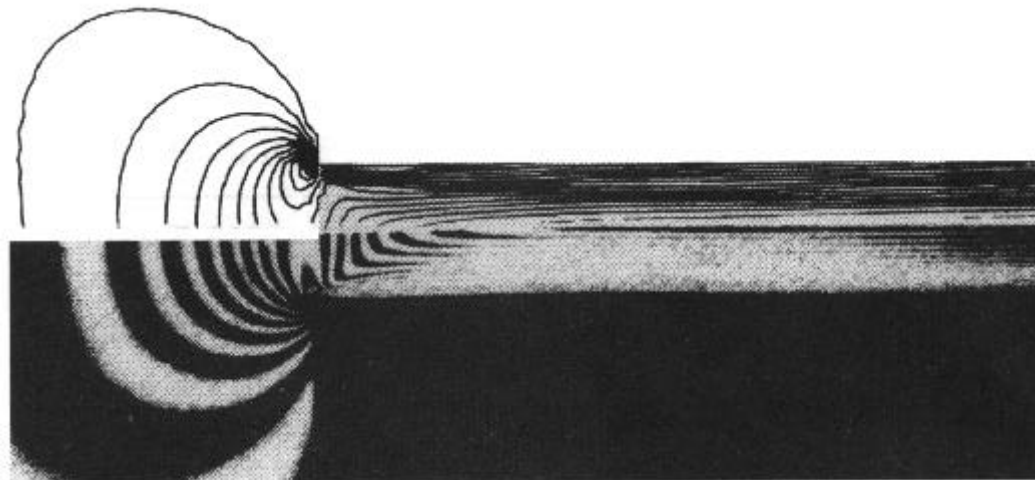
- ⊆ Normal stress ratio formed with rheological data from homogeneous transient uniaxial extension and simple steady shear flow.

...but entry flow is a complex flow with nonhomogeneous extensional kinematics

- ⊆ Would be nice to be able to form normal stress ratio from *in situ* stress measurements.

Flow Induced Birefringence (FIB) $\propto \gamma \dot{\gamma} n \mathbb{N}, P \propto (J_{zz} - J_{rr}), J_{rz}$

- ⊆ FIB measurements are also an excellent comparative tool for evaluating the ability of constitutive models to capture small scale physics.



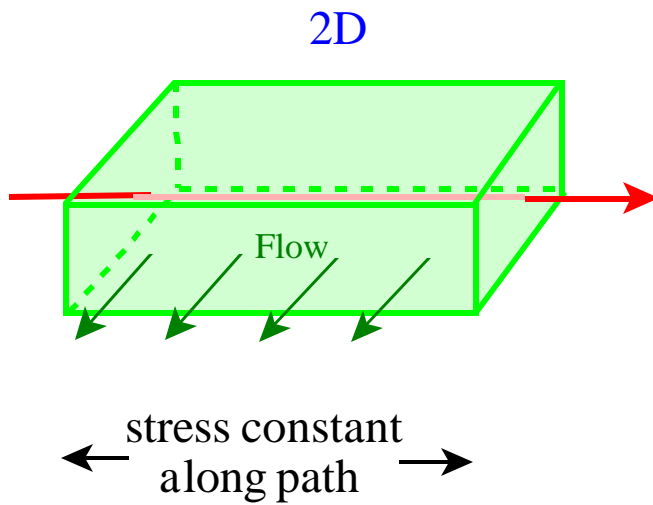
Bernaudo et al. (1998) - Flow of a LLDPE melt into a 8:1 planar contraction

Axisymmetric Flow Induced Birefringence



Li and Burghardt (1995)

Ⓒ FIB typically used in two-dimensional flows because it is a line-of-sight technique.



birefringence stress-optical coefficient

$$\Delta n' \sin 2\chi = 2C n_m k_B T ?_{12}$$

$$\Delta n' \cos 2\chi = C n_m k_B T (?_{11} - ?_{22})$$

extinction angle

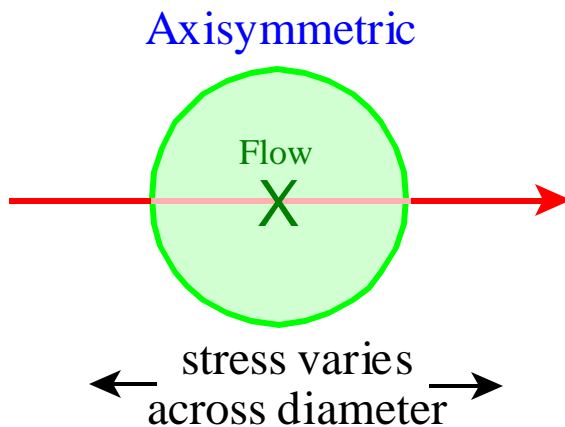
$$\delta = \frac{2\pi \Delta n' d}{\lambda}$$

retardation

$$\mathbf{A} = \langle \mathbf{Q}\mathbf{Q} \rangle$$

chain end-to-end vector

Ⓒ Axisymmetric flow result in an integrated measure of FIB and polymer chain conformation.

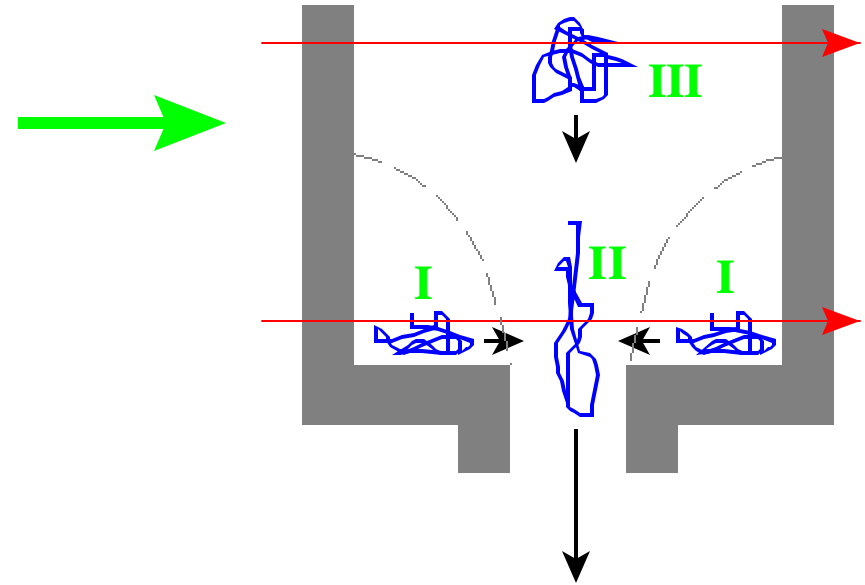
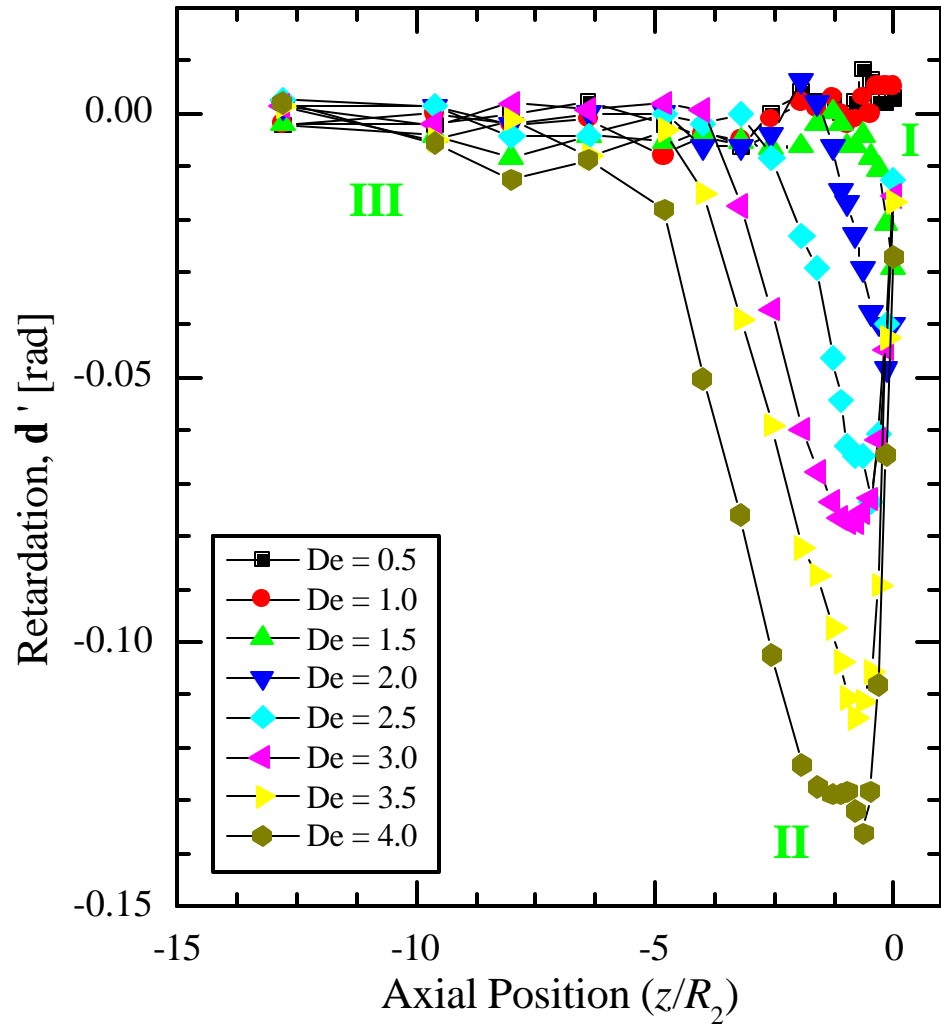


Ⓒ Integrated measures dependent on optical train used.

$$I_1 = \sin \delta' \cos 2\chi' = \int \sin \delta \cos 2\chi \, dl$$

$$I_2 = [1 - \cos \delta'] \sin 2\chi' \cos 2\chi' = \int [1 - \cos \delta] \sin 2\chi \cos 2\chi \, dl$$

Axisymmetric FIB Upstream of 4:1:4 Contraction-Expansion

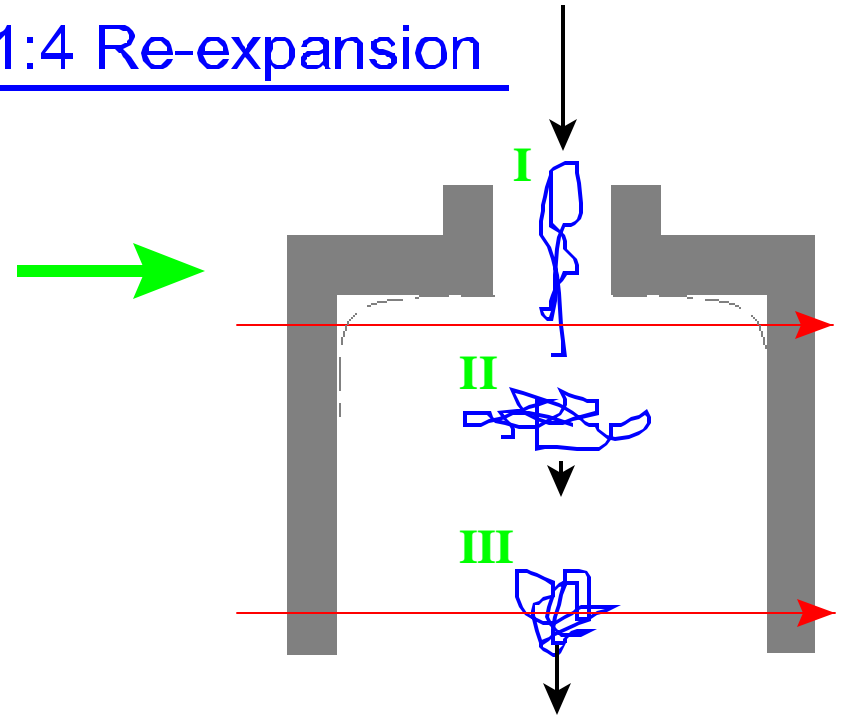
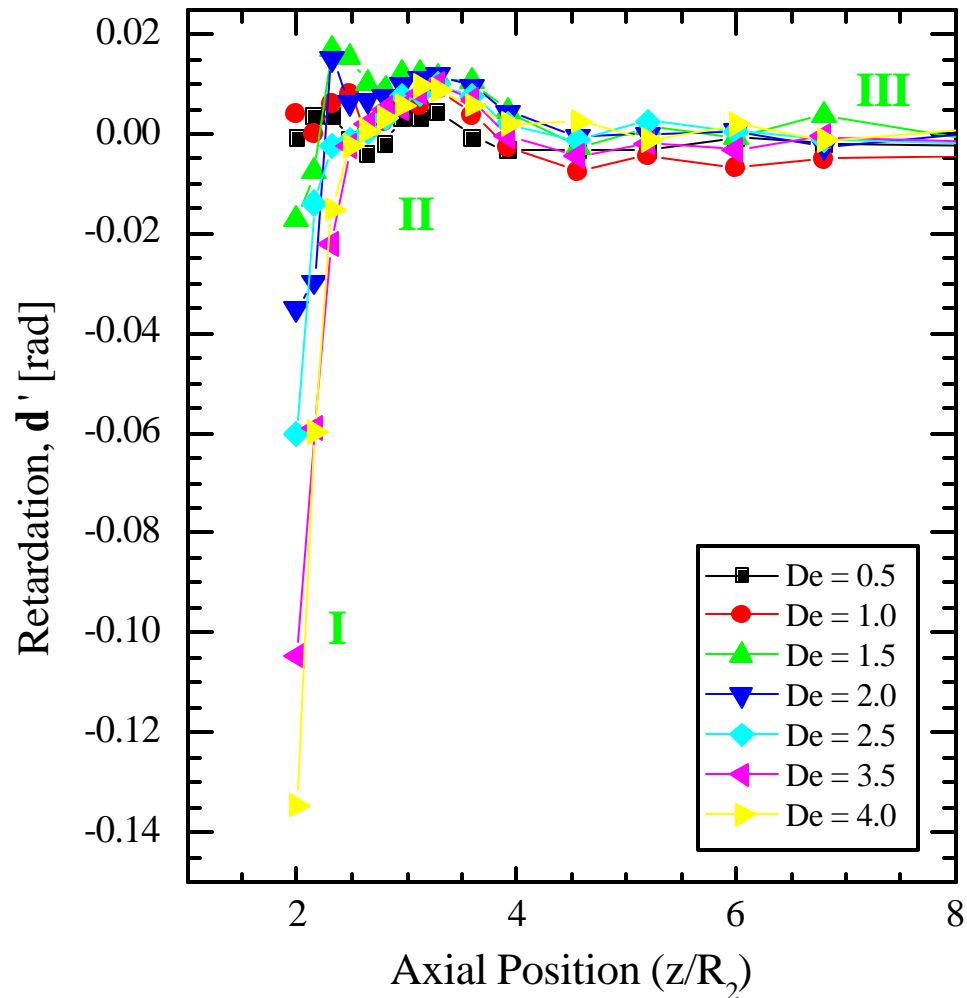


C Axial compression ($d\| > 0$) near contraction ($z = 0$) due to shear flow along contraction plane.

C At large De , strong axial elongation ($d\| < 0$) results from extensional flow along centerline.

C Growth in size and strength of axial elongation region coincides with corner vortex growth.

Axisymmetric FIB Downstream of 4:1:4 Re-expansion

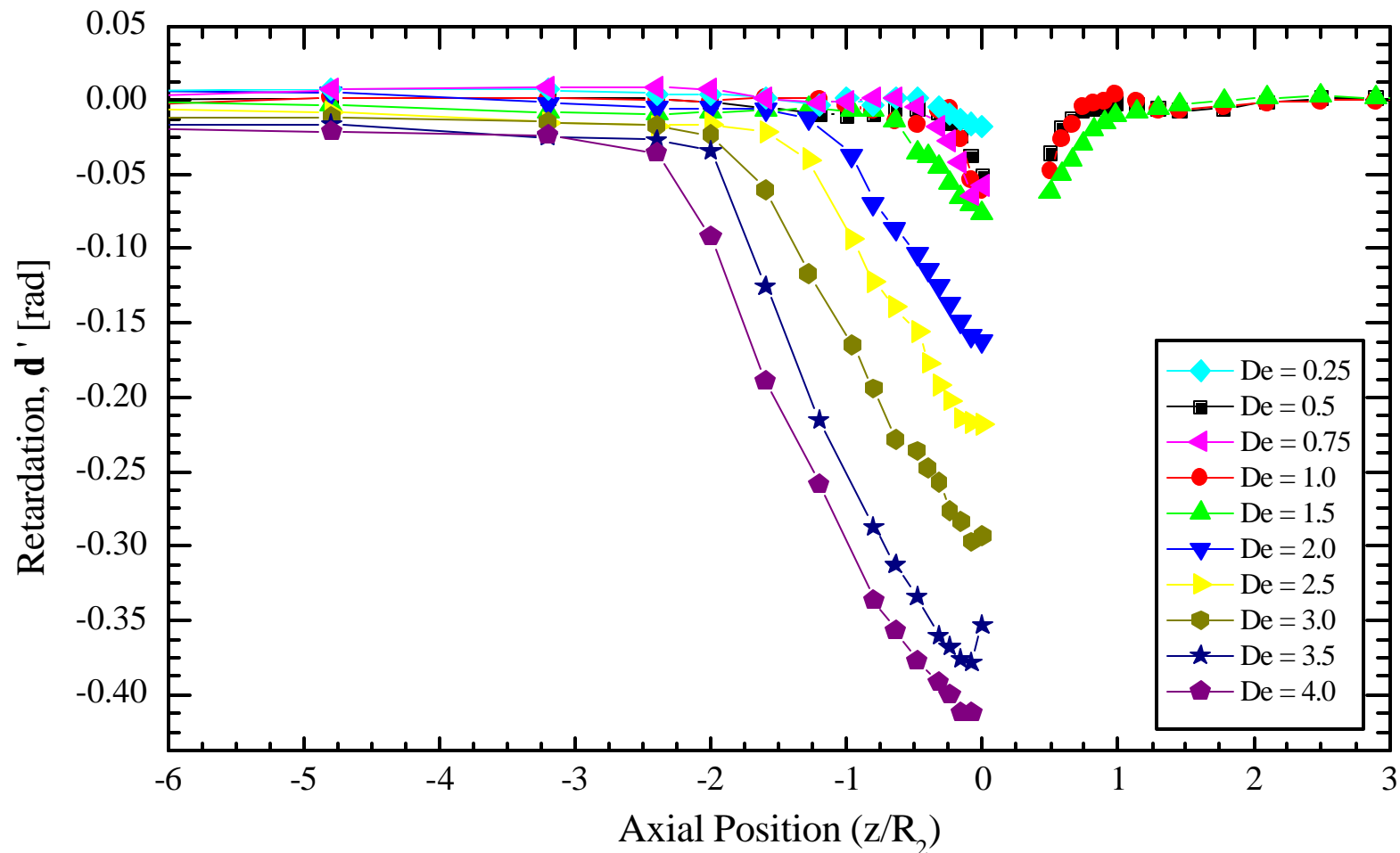


C Region of axial compression ($z=-3$) resulting from deceleration of flow out of contraction.

C Magnitude and amplitude of compression doesn't change with Deborah number.

C At large Deborah numbers, strong axial elongation near re-expansion plane resulting from extensional flow upstream and strong shear in the throat.

Axisymmetric FIB in a 2:1:2 Contraction-Expansion



- C Again, strong extension is observed upstream corresponding to vortex growth.
- C Very little compression is observed from the shearing flow at the contraction plane.
- Y No other distinct qualitative differences are observed between different contraction ratios.

Normal Stress Ratio from FIB Measurements



Can now form a normal stress ratio from axisymmetric FIB measurements.

$$\aleph = \frac{N_1 / \eta_0 \dot{\gamma}}{(\tau_{zz} - \tau_{rr}) / \eta_0 \dot{\epsilon}} = \frac{S_R(\dot{\gamma})}{T_R(\dot{\epsilon})}$$

To calculate normal stress ratio from FIB measurements we:

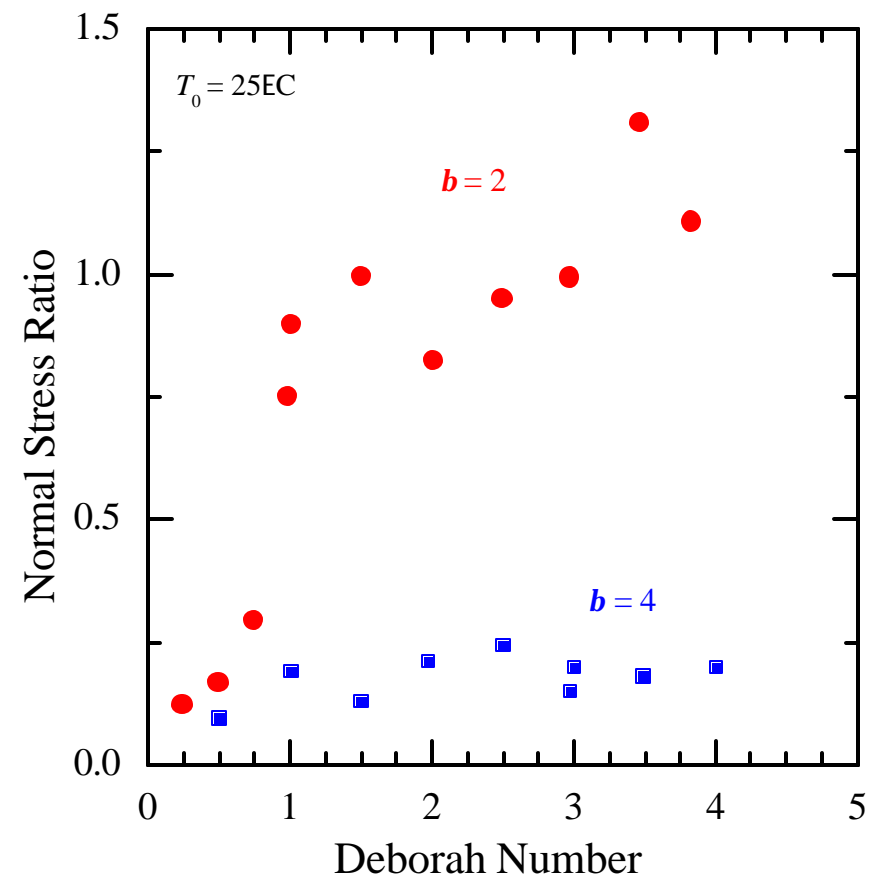
6 Use simple shear rheology to determine N_1 at the contraction wall.

6 Assume all extensional stress is held by a single uniformly stretched filament of width $2R_2$.

$$(\tau_{zz} - \tau_{rr}) = \delta\lambda / 4\pi CR_2$$

Normal stress ratio strongly dependent on contraction ratio.

Final test would be to compare with FIB measurements of PIB/PB Boger fluids.



Conclusions



- ③ We have coupled global flow field measurements of the effect of viscoelasticity with local conformation measurements for a model polymer solution in a prototypical complex flow.
- ③ These measurements generate a comprehensive data set for the validation of constitutive equations and numerical methods.
 - ⑥ Pressure drop measurements
 - ⑥ Axisymmetric flow induced birefringence measurements
 - ⑥ Velocity measurements (PIV and LDV)
 - ⑥ Vortex growth measurements
 - ⑥ Characterization of an elastic flow instability
- ③ We have rationalized the dependence of elastic lip and corner vortices on contraction ratio and test fluid rheology with a new dimensionless group, the normal stress ratio.

$$\mathfrak{N} = \frac{N_1 / \eta_0 \dot{\gamma}}{(\tau_{zz} - \tau_{rr}) / \eta_0 \dot{\epsilon}} = \frac{S_R(\dot{\gamma})}{T_R(\epsilon)}$$

Flow Induced Birefringence (FIB)



C Can use FIB to determine microscopic anisotropy in polymer chain conformation

birefringence

stress-optical coefficient

chain end-to-end vector

$$\Delta n' \sin 2\chi = 2C n_m k_B T \langle Q^2 \rangle$$

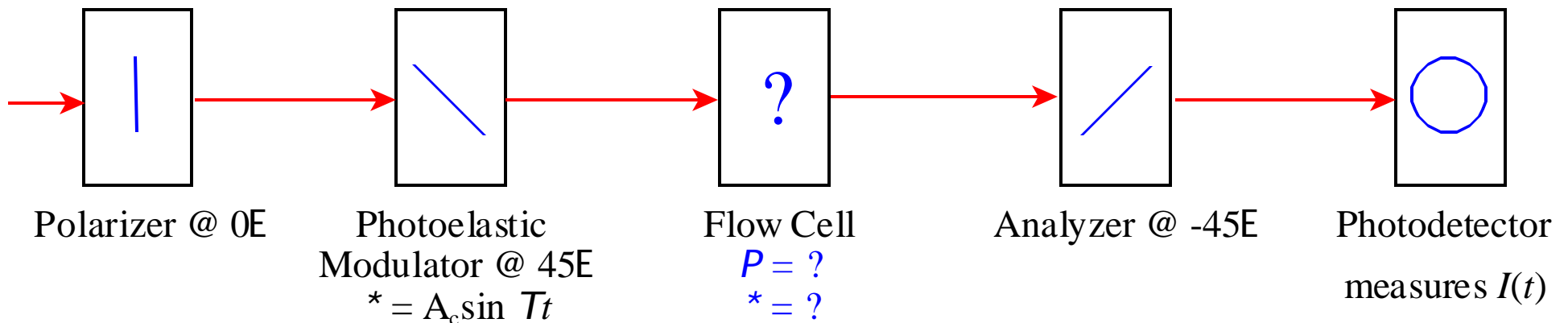
$$\mathbf{A} = \langle \mathbf{Q}\mathbf{Q} \rangle$$

$$\Delta n' \cos 2\chi = C n_m k_B T (\langle Q_{11} \rangle - \langle Q_{22} \rangle)$$

$$\delta = 2\pi \Delta n' d / \lambda$$

C Need a technique to measure retardance ($*$) and extinction angle (P) simultaneously

*** Polarization Modulated Flow Birefringence ***



$$I(t) = I_{dc} [1 + 2J_1(A_c) M_{34} \sin \omega t + 2J_2(A_c) M_{32} \cos 2\omega t]$$

$$M_{34} = \sin \delta \cos 2\chi$$

$$M_{32} = [1 - \cos \delta] \sin 2\chi \cos 2\chi$$



$$R_\omega = \frac{I_\omega}{2J_1(A_c) I_{dc}} = M_{34},$$

$$R_{2\omega} = \frac{I_{2\omega}}{2J_2(A_c) I_{dc}} = M_{32},$$